# Singular perturbations of gradient flows and rate-independent evolution 

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CALCULUS OF VARIATIONS AND APPLICATIONS
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In collaboration with Virginia Agostiniani, Riccarda Rossi

## Outline

1 Rate-independent evolution and singular perturbation of gradient flows

2 Transversality conditions for the critical set

3 Compactness and variational characterization of the limit evolution

4 A useful tool: graph convergence

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## Evolution by critical/stable points

- $\mathbb{H}:=\mathbb{R}^{\mathrm{d}}(\rightsquigarrow$ Hilbert space),
- $\mathcal{E}:[0, \mathrm{~T}] \times \mathbb{H} \rightarrow \mathbb{R}$ is a $\mathrm{C}^{2}$ time dependent energy with $\mathbb{H}$-differential $\mathrm{D} \mathcal{E}:[0, \mathrm{~T}] \times \mathbb{H} \rightarrow \mathbb{H}$.
Typical example: time dependent linear perturbation

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- $\mathbf{u}_{0} \in \mathbb{H}$.
- Critical points:

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\mathbf{C} & :=\{(t, x): D \mathcal{E}(\mathrm{t}, \mathrm{x})=0\}, \\
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$$

- $\rho$-critical points: fix some $\rho>0 \quad\|\mathrm{D}(\mathrm{t}, \mathrm{x})\| \leqslant \rho$
- Globally $\rho$-stable points:

$$
\mathcal{E}(\mathrm{t}, \mathrm{x}) \leqslant \mathcal{E}(\mathrm{t}, \mathrm{y})+\rho\|y-x\| \quad \forall y \in \mathbb{H}
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Globally $\rho$-stable points are $\rho$-critical.

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## Aim:

Select "reasonable" evolution curves $t \mapsto \mathbf{u}(t)$ starting from $\mathbf{u}_{0}$ such that $\mathbf{u}(t)$ is critical/stable for every time $t \in[0, T]$.

Simple examples in 1D



Double well





## Time Incremental Minimization Scheme

In the case of global $\rho$-stable evolutions, the main tool to provide existence and to approximate solutions is

The time Incremental Minimization scheme
Fix $\tau:=T / N$ (for simplicity), $t_{\tau}^{n}:=n \tau, U_{\tau}^{0}=\mathbf{u}(0)$.
Recursively choose $\mathrm{U}_{\tau}^{n}$ among the minimizers of

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\mathrm{U} \mapsto \mathcal{E}\left(\mathrm{t}_{\tau}^{\mathrm{n}}, \mathrm{U}\right)+\rho\left\|\mathrm{U}-\mathrm{U}_{\tau}^{\mathrm{n}-1}\right\|
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$\overline{\mathrm{U}}_{\tau}:=$ the piecewise constant interpolant of the values $\mathrm{U}_{\tau}^{n}$.

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$\overline{\mathrm{U}}_{\tau}:=$ the piecewise constant interpolant of the values $\mathrm{U}_{\tau}^{n}$.
Theorem [Mainik-Mielke '05]:
There exists a sequence $\mathrm{k} \mapsto \tau(\mathrm{k}) \downarrow 0$ and $\mathrm{u}:[0, \mathrm{~T}] \rightarrow \mathbb{H}$ such that and

$$
\overline{\mathrm{U}}_{\tau(\mathrm{k})}(\mathrm{t}) \rightarrow \mathbf{u}(\mathrm{t}), \quad \mathcal{E}\left(\mathrm{t}, \overline{\mathrm{U}}_{\tau(\mathrm{k})}(\mathrm{t}) \rightarrow \mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t})) \quad \text { for every } \mathrm{t} \in[0, \mathrm{~T}]\right.
$$

and $\mathbf{u}$ is called an Energetic solution to the Rate Independent Ssystem (R.I.S.) $(\mathbb{H}, \mathcal{E}, \rho)$.

## Energetic solutions

Energetic solution: a curve $u:[0, T] \rightarrow \mathbb{H}$ satisfying for every $t \in[0, T]$ the $\rho$-stability condition

$$
\begin{equation*}
\varepsilon(\mathrm{t}, \mathbf{u}(\mathrm{t})) \leqslant \varepsilon(\mathrm{t}, v)+\rho\|\mathbf{u}(\mathrm{t})-v\| \quad \text { for every } v \in \mathbb{H}, \tag{S}
\end{equation*}
$$

and the energy balance

$$
\begin{equation*}
\mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t}))+\rho \operatorname{Var}(\mathbf{u},[0, \mathrm{t}])=\mathcal{E}(0, \mathbf{u}(0))+\int_{0}^{\mathrm{t}} \mathcal{P}(\mathrm{r}, \mathbf{u}(\mathrm{r})) \mathrm{dr} \tag{E}
\end{equation*}
$$

where

$$
\mathcal{P}(t, x)=\frac{\partial}{\partial t} \mathcal{E}(t, x) .
$$

[Mielke-Theil-Levitas '02, Mielke-Theil '04
Francfort-Marigo '93/'98, DalMaso-Toader '02, Francfort-Larsen '05,
DalMaso-Francfort-Toader '05
Mainik-Mielke '05, Francfort-Mielke '06
Mielke-Roubicek '15]

## The "smooth" finite dimensional case

Energetic solutions provides a variational selection among trajectories satisfying

$$
\rho \operatorname{sign}(\dot{\mathbf{u}}(\mathrm{t}))+\mathrm{D} \mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t})) \ni 0 \quad \text { in particular }\|\mathrm{D} \mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t}))\| \leqslant \rho,
$$

and at every jump point $t \in J(u)$ the energetic jump conditions

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\mathcal{E}(\mathrm{t},(\mathbf{u}(\mathrm{t}-))-\mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t}+))=\rho\|\mathbf{u}(\mathrm{t}+)-\mathbf{u}(\mathrm{t}+)\|
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Energetic solution in the 1-dimensional case for a strictly increasing f.

## A few technical points

- Compactness w.r.t. space: it follows by the compactness of the sublevels of $E$ and the estimate

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- Compactness w.r.t. time: it follows by the uniform BV estimate (Helly's Theorem)

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- Stability: it follows by the stability of each minimizer and from the

$$
\text { closure of the } \rho \text {-stable set }\{(t, x): E(t, x) \leqslant E(t, y)+\rho\|y-x\|\}
$$

## Viscous corrections of the Incremental Minimization Scheme

By introducing a small parameter $\varepsilon=\varepsilon(\tau)>0$, one may consider the following modified incremental minimization scheme

$$
\text { minimize } \quad \mathrm{U} \mapsto \mathcal{E}\left(\mathrm{t}_{\tau}^{\mathrm{n}}, \mathrm{U}\right)+\rho\left\|\mathrm{U}-\mathrm{U}_{\tau}^{\mathrm{n}-1}\right\|+\frac{\varepsilon}{2 \tau}\left\|\mathrm{U}-\mathrm{U}^{\mathrm{n}-1}\right\|^{2}
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and its limit behaviour in three cases:

- Visco Energetic solutions: [Minotti-S.]

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\varepsilon / \tau=\mu>0 \text { and } \rho>0 \text { are fixed. } \tag{VE}
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The last two methods correspond to the limit behaviour of

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\rho \operatorname{sign}(\dot{\mathbf{u}}(\mathrm{t}))+\varepsilon \dot{\mathbf{u}}(\mathrm{t})+\mathrm{D} \mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t})) \ni 0
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## Singular limit of gradient flows $(\rho=0)$

## Main problem

Study the asymptotic behaviour as $\varepsilon \downarrow 0$ of the solution $\mathfrak{u}_{\varepsilon}:[0, \mathrm{~T}] \rightarrow \mathbb{H}$ of the gradient flow

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\left\{\begin{aligned}
\varepsilon \mathbf{u}_{\varepsilon}^{\prime}(\mathrm{t}) & =-\mathrm{D} \mathcal{E}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right) \\
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In order to avoid a transition layer at $t=0$ we will assume $D \mathcal{E}\left(0, \mathbf{u}_{0}\right)=0$.

## Basic energy estimate

Set $\mathcal{P}(t, x)=\partial_{t} \mathcal{E}(t, x)$. Chain rule: $\frac{d}{d t} \mathcal{E}(t, x(t))=\left\langle D \mathcal{E}(t, x), x^{\prime}(t)\right\rangle+\mathcal{P}(t, x(t))$.

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## Energy identity

$\mathcal{E}\left(T, \mathbf{u}_{\varepsilon}(\mathrm{T})\right)+\int_{0}^{T}\left(\frac{\varepsilon}{2}\left\|\mathbf{u}_{\varepsilon}^{\prime}(\mathrm{t})\right\|^{2}+\frac{1}{2 \varepsilon}\left\|D \mathcal{E}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right)\right\|^{2}\right) d t=\mathcal{E}\left(0, \mathbf{u}_{0}\right)+\int_{0}^{T} \mathcal{P}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right) \mathrm{dt}$

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$$

## Basic energy estimate

Set $\mathcal{P}(\mathrm{t}, \mathrm{x})=\partial_{\mathrm{t}} \mathcal{E}(\mathrm{t}, \mathrm{x})$. Chain rule: $\frac{\mathrm{d}}{\mathrm{dt}} \mathcal{E}(\mathrm{t}, \mathrm{x}(\mathrm{t}))=\left\langle\mathrm{D} \mathcal{E}(\mathrm{t}, \mathrm{x}), \mathrm{x}^{\prime}(\mathrm{t})\right\rangle+\mathcal{P}(\mathrm{t}, \mathrm{x}(\mathrm{t}))$.

$$
\varepsilon\left\|\mathbf{u}_{\varepsilon}^{\prime}(\mathrm{t})\right\|^{2}=\frac{1}{\varepsilon}\left\|\mathrm{D} \mathcal{E}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right)\right\|^{2}=-\frac{\mathrm{d}}{\mathrm{dt}} \mathcal{E}(\mathrm{t}, \mathrm{x}(\mathrm{t}))+\mathcal{P}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right)
$$

## Energy identity

$\mathcal{E}\left(T, \mathbf{u}_{\varepsilon}(\mathrm{T})\right)+\int_{0}^{T}\left(\frac{\varepsilon}{2}\left\|\mathbf{u}_{\varepsilon}^{\prime}(\mathrm{t})\right\|^{2}+\frac{1}{2 \varepsilon}\left\|\mathrm{D} \mathcal{E}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right)\right\|^{2}\right) \mathrm{dt}=\mathcal{E}\left(0, \mathbf{u}_{0}\right)+\int_{0}^{T} \mathcal{P}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right) \mathrm{dt}$
If

$$
\lim _{|x| \rightarrow \infty} \mathcal{E}(t, x)=+\infty, \quad \mathcal{P}(t, x) \leqslant A+B|\mathcal{E}(t, x)|
$$

## Basic estimates

$$
\begin{aligned}
\left\|\mathbf{u}_{\varepsilon}(\mathrm{t})\right\| & \leqslant \mathrm{C} \\
\mathcal{E}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right) & \leqslant \mathrm{C} \\
\int_{0}^{T}\left\|\mathbf{u}_{\varepsilon}^{\prime}(\mathrm{t})\right\|^{2} \mathrm{dt} & \leqslant \mathrm{C} / \varepsilon \\
\int_{0}^{T} \| \mathrm{D} \mathcal{E}\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t}) \|^{2} \mathrm{dt}\right. & \leqslant \mathrm{C} \varepsilon
\end{aligned}
$$

## Main difficulties

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When $\mathcal{E}$ is not convex, one may expect jumps and more complex bifurcation behaviour when $\mathbf{u}(t)$ hits a degenerate critical point of

$$
\mathbf{C}_{\mathrm{d}}:=\left\{(\mathrm{t}, \mathrm{x}) \in \mathbf{C}(\mathrm{t}): \mathrm{D}_{x x}^{2} \mathcal{E}(\mathrm{t}, \mathrm{x}) \quad \text { is not invertible }\right\}
$$

The critical set (1D)


Jumps between curves


## Outline

1 Rate-independent evolution and singular perturbation of gradient flows

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## Strong transversality [Zanini, '08]

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In 1 D the above conditions mean that $\mathrm{G}(\mathrm{t}, \mathrm{x}):=\partial_{x} \mathcal{E}(\mathrm{t}, \mathrm{x})$ satisfies
whenever $G(t, x)=0, \partial_{x} G(t, x)=0$ then $\partial_{t} G(t, x) \neq 0, \quad \partial_{x x}^{2} G(t, x) \neq 0$.
Strong transversality implies that $\mathbf{C}(\mathrm{t})$ is discrete for every $\mathrm{t} \in[0, \mathrm{~T}]$. [Agostiniani-Rossi, Scilla-Solombrino].

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It is a "generic" condition, in the following sense: if $\mathcal{E}$ is $\mathrm{C}^{4}$, for a $\mathrm{G}_{\delta}$ dense set of $\mathrm{g} \in \mathbb{H}$ and $\mathrm{Q} \in \mathscr{L}(\mathbb{H}, \mathbb{H})$ the perturbed energy

$$
\mathcal{E}_{\mathrm{g}, \mathrm{Q}}(\mathrm{t}, \mathrm{x}):=\mathcal{E}(\mathrm{t}, \mathrm{x})-\langle\mathrm{g}, \mathrm{x}\rangle-\langle\mathrm{Qx}, \mathrm{x}\rangle
$$

satisfies the strong transversality conditions.

## Transversality of the critical set



## Simpler Generic conditions

We consider only linear perturbations

$$
\mathcal{E}_{\mathrm{g}}(\mathrm{t}, \mathrm{x}):=\mathcal{E}(\mathrm{t}, \mathrm{x})-\langle\mathrm{g}, \mathrm{x}\rangle, \quad \mathrm{D} \mathcal{E}_{\mathrm{g}}(\mathrm{t}, \mathrm{x})=\mathrm{D} \mathcal{E}(\mathrm{t}, \mathrm{x})-\mathrm{g} .
$$

A generic decomposition of C [Sard, Quinn, Hirsch, Simon; Saut-Temam]
The set $\mathscr{O}$ of $g \in \mathbb{H}$ such that the total differential

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\mathrm{dD} \mathcal{E}_{\mathrm{g}}(\mathrm{t}, \mathrm{x}) \in \mathscr{L}(\mathbb{R} \times \mathbb{H} ; \mathbb{H}) \quad \text { is surjective for every }(\mathrm{t}, \mathrm{x}) \in \mathbf{C}
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If $\mathbf{g} \in \mathscr{O}, \mathbf{C}$ can be decomposed as the disjoint union of at most countable $\mathrm{C}^{1}$ curves. In particular

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- $\mathcal{E}(\mathrm{t}, \cdot)$ is constant on every connected component of $\mathbf{C}(\mathrm{t})$.
- There is an at most countable set of times $N \subset[0, T]$ such that $\mathbf{C}(t)$ is not totally disconnected:
- if $t \in[0, T] \backslash N$ then every connected component of $\mathbf{C}(t)$ is reduced to a point,
- if $t \in N$ then $\mathbf{C}(t)$ has "larger" connected components.


## A simple example in infinite dimension

Let $\Omega$ be a bounded connected open set of $\mathbb{R}^{3}, \mathbb{H}:=\mathrm{L}^{2}(\Omega)$,
$D(E):=H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. Consider

$$
\begin{aligned}
E(u) & :=\int_{\Omega}\left(\frac{1}{2}|\nabla u(x)|^{2}+W(u(x))\right) d x, \quad f \in C^{2}\left([0, T] ; L^{2}(\Omega)\right) \\
\mathcal{E}(t, u) & =E(u)-\int_{\Omega} f(t, x) u(x) d x, \quad \text { if } u \in D(E) . \\
-\Delta u(t, x) & +W^{\prime}(u(t, x))-g(x)=f(t, x) \quad \text { in } \Omega, \quad u(t, \cdot)=0 \text { on } \partial \Omega .
\end{aligned}
$$

For a dense $\mathrm{G}_{\delta}$-subset $\mathscr{O}$ in $\mathrm{L}^{2}(\Omega)$ the energy $\mathcal{E}_{g}, \mathrm{~g} \in \mathscr{O}$, satisfies the transversality condition and the critical set is countably $\left(\mathscr{H}^{1}, 1\right)$-rectifiable.

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One can also consider genericity w.r.t. $\Omega$ or w.r.t. the coefficients of the elliptic operator.

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## The main compactness result

## Theorem (Agostiniani-Rossi-S.)

Suppose that $\mathbf{C}$ is countably $\left(\mathscr{H}^{1}, 1\right)$-rectifiable (it is sufficient that $\mathscr{H}^{1}$ is $\sigma$-finite on C)

Then there exists:

- a subsequence $\mathrm{n} \mapsto \varepsilon(\mathrm{n}) \downarrow 0$
- a limit curve $\mathbf{u}:[0, \mathrm{~T}] \rightarrow \mathbb{H}$ such that

$$
\lim _{n \rightarrow \infty} \mathbf{u}_{\varepsilon(n)}(t)=\mathbf{u}(t) \quad \text { for every } t \in[0, T]
$$

## Properties of limit solutions

Let $\mathfrak{u}:[0, \mathrm{~T}] \rightarrow \mathbb{H}$ be a limit solution arising from the previous compactness result.

For the sake of simplicity, we suppose that $\mathbf{C}(\mathrm{t})$ is totally disconnected for every $t \in[0, T]$, i.e. $N=\emptyset$.

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- $u$ is regulated, i.e. for every $t \in[0, T]$
$\exists \lim _{s \uparrow t} \mathbf{u}(\mathrm{~s})=\mathbf{u}(\mathrm{t}-)$,
$\exists \lim _{s \downarrow t} \mathbf{u}(\mathrm{~s})=\mathbf{u}(\mathrm{t}+)$,
$\mathrm{J}(\mathbf{u}):=\{\mathrm{t} \in[0, \mathrm{~T}]: \mathbf{u}(\mathrm{t}) \neq \mathbf{u}(\mathrm{t} \pm)\}$ is at most countable.


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- $\mathbf{u}(\mathrm{t}) \in \mathbf{C}(\mathrm{t})$ for every $\mathrm{t} \in[0, \mathrm{~T}] \backslash \mathrm{J}(\mathbf{u})$; if $\mathrm{t} \in \mathrm{J}(\mathrm{u})$ then $\mathbf{u}(\mathrm{t}-), \mathbf{u}(\mathrm{t}+) \in \mathbf{C}(\mathrm{t})$.


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- The map $t \mapsto \mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t}))$ has bounded variation, its jump set coincides with the jump set of $u$ for every $0 \leqslant s<t \leqslant T$ the energy balance holds:

$$
\mathcal{E}(t, \mathbf{u}(t-))+\sum_{r \in J(u) \cap(s, t)} c(r ; \mathbf{u}(r-), \mathbf{u}(r+))=\mathcal{E}(s, \mathbf{u}(s+))+\int_{s}^{t} \mathcal{P}(r, \mathbf{u}(r)) d r
$$

## The transition cost c

At every jump $t \in J(u)$, the energy dissipation corresponds to the optimal transition cost:

$$
\mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t}-))-\mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t}+))=\mathrm{c}(\mathrm{t} ; \mathbf{u}(\mathrm{t}-), \mathbf{u}(\mathrm{t}+))
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For every $\mathrm{t} \in[0, \mathrm{~T}]$ the transition cost $\mathrm{c}\left(\mathrm{t} ; \mathrm{u}_{-}, \mathrm{u}_{+}\right)$between two points $u_{-}, u_{+} \in \mathbb{H}$ is given by the "Finsler" metric induced by $\|D \mathcal{E}\|$ :

$$
c\left(t ; u_{-}, u_{+}\right):=\inf \left\{\int_{\Omega(\vartheta)}\|D \mathcal{E}(t, \vartheta(\tau))\| \| \vartheta^{\prime}(\tau)\right) \| d \tau:
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& \left.\vartheta \in \operatorname{Lip}_{\mathrm{loc}}(\Omega), \mathcal{E}(\mathrm{t}, \vartheta(\cdot)) \in \operatorname{Lip}([0,1])\right\}
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& \left.\vartheta \in \operatorname{Lip}_{\mathrm{loc}}(\Omega), \mathcal{E}(\mathrm{t}, \vartheta(\cdot)) \in \operatorname{Lip}([0,1])\right\}
\end{aligned}
$$

We always have

$$
\mathcal{E}\left(\mathrm{t}, \mathbf{u}_{-}\right)-\mathcal{E}\left(\mathrm{t}, \mathbf{u}_{+}\right) \leqslant \mathrm{c}\left(\mathrm{t} ; \mathbf{u}_{-}, \mathbf{u}_{+}\right)
$$

Transitions


$$
C(t)=\left\{x_{1}, x_{2}, x_{3}, x_{L}\right\} \quad \int_{z_{-}}^{\tau_{t}}|D \varepsilon(t, \theta(t))| \cdot\left|\theta^{\prime}(t)\right| d \tau-\operatorname{simin}
$$

## Structure of limit solutions

The limit energy identity

$$
\mathcal{E}(\mathrm{t}, \mathbf{u}(\mathrm{t}-))+\sum_{\mathrm{r} \in \mathrm{~J}(\mathbf{u}) \cap(\mathrm{s}, \mathrm{t})} \mathbf{c}(\mathrm{r} ; \mathbf{u}(\mathrm{r}-), \mathbf{u}(\mathrm{r}+))=\mathcal{E}(\mathrm{s}, \mathbf{u}(\mathrm{~s}+))+\int_{\mathrm{s}}^{\mathrm{t}} \mathcal{P}(\mathrm{r}, \mathbf{u}(\mathrm{r})) \mathrm{dr}
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- in each connected component of $\Omega(\vartheta) \subset[0,1]$ where $\vartheta \notin \mathbf{C}(\mathrm{t})$ there is an increasing change of variable $\tau=\tau(r), r \in(a, b)$, such that the reparametrized transition $\theta(\mathrm{r}):=\vartheta(\tau(\mathrm{r}))$ satisfies the gradient flow equation

$$
\frac{\mathrm{d}}{\mathrm{dr}} \theta(\mathrm{r})=-\mathrm{D} \mathcal{E}(\mathrm{t}, \theta(\mathrm{r}))
$$

at the "frozen time" t .

## Jumps in the smooth 1-dimensional case: double-well potential



Limit solution in the 1-dimensional case for a strictly increasing f. The blue line represents the graph of the jump transition $\vartheta$.

## Outline

1 Rate-independent evolution and singular perturbation of gradient flows

2 Transversality conditions for the critical set

3 Compactness and variational characterization of the limit evolution

4 A useful tool: graph convergence

## Main idea

Instead of studying the convergence of $u_{\varepsilon}$ we consider the limit of their graphs:

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\mathbf{G}_{\varepsilon}:=\left\{\left(\mathrm{t}, \mathbf{u}_{\varepsilon}(\mathrm{t})\right): \mathrm{t} \in[0, \mathrm{~T}]\right\} \subset[0, \mathrm{~T}] \times \mathbb{H}
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For every $\varepsilon>0$ there exists $\bar{n} \in \mathbb{N}$ such that for every $n \geqslant \bar{n}$
$\mathrm{K}_{\mathrm{n}}$ is contained in the $\varepsilon$-neighborhood of K ,
$K$ is contained in the $\varepsilon$-neighborhood of $K_{n}$,

## A general compactness property

## Blashke compactness theorem

If $\mathbf{G}_{\varepsilon}$ are contained in a common compact set, there exists a subsequence $n \mapsto \varepsilon(n) \downarrow 0$ and a limit set $\mathbf{G}$ such that $\mathbf{G}_{\varepsilon(n)} \xrightarrow{K} \mathbf{G}$.

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If moreover the sets $\mathbf{G}_{\varepsilon}$ are connected then also the limit $\mathbf{G}$ is connected.

## Examples



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$N M N$ AMMA.

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## Extensions

- The assumption that $\mathbf{C}(t)$ is totally disconnected for every $t \in[0, T]$ can be removed, by working with possibly discontinuous functions at a countable set of times $N \subset[0, T]$. For each $t \in N$ we can only say that $u(s)$ approaches a connected component of $\mathbf{C}(\mathrm{t})$ as $\mathrm{s} \rightarrow \mathrm{t}$.


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- A finer description is possible when $\mathcal{E}$ is subanalytic.

