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R^n gravity is kicking and alive: the cases of Orion and NGC 3198

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We analyzed the Rotation Curves of two crucial objects, the Dwarf galaxy Orion and the low luminosity Spiral NGC 3198, in the framework of R^n gravity. We surprisingly found that the no DM power-law $F(R)$ case fits them well, performing much better than LCDM Dark Matter halo models. The level of this unexpected success can be a boost for R^n gravity.

Keywords: Dark matter; R^n gravity; alternative theories of gravity.

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1. Introduction

It is well-known that the Rotation Curves (RCs) of spiral galaxies show a non-Keplerian circular velocity profile which cannot be explained by considering a Newtonian gravitational potential generated by the baryonic matter.¹⁵ Current possible explanations include the postulate of a new yet not detected state of matter, the dark matter, e.g.,¹⁵ a phenomenological modification of the Newtonian dynamics,¹² and higher order Gravitational Theories, see e.g.^{2,4,6,9-11}

A recent theory proposed by⁵ modifies the usual Newtonian gravitational potential generated by (baryonic) matter as an effect of power-law fourth order theories of gravity that replace in the gravity action the Ricci scalar R with a function $f(R) \propto R^n$, where n is a slope parameter. The goal is that the galaxy kinematics resulting in the $f(R)$ scenario from the luminous matter alone would account for those observations that the front runner candidate of the competing scenario, i.e. a Cold Dark Matter particle, fails to account.

In the current theory the Newtonian potential generated by a point-like source gets modified in to

$$\phi(r) = -\frac{Gm}{r} \left\{ 1 + \frac{1}{2} [(r/r_c)^\beta - 1] \right\}, \quad (1)$$

where β is a function of the slope n , and r_c is a scale length parameter. At a fixed n , β is a universal constant, while r_c depends on the particular gravitating system being studied. In a virialized system the circular velocity is related to the derivative of the potential through $V^2 = r d\phi(r)/dr$. In short, can Eq. (1) explain, without a Dark Component, the circular velocity in spirals and specially that in cases in which halos of (Cold) Dark Particles fail?

Frigerio Martins and Salucci⁸ investigated the consistency and the universality of this theory by means of a sample of spirals, obtaining a quite good success that was encouraging for further investigations. Recently, crucial information for two special objects has been available and we are now able to test the theory in unprecedented accurate way.

Orion is a dwarf galaxy of luminous mass $< \frac{1}{100}$ the Milky Way stellar disk mass with a baryonic distribution dominated by a HI disk, whose surface density is accurately measured and, noticeably, found to have some distinct feature. The stellar disk, on the other side, is a pure exponential disk. The available rotation curve¹⁷ is extended and it is of very high resolution. Noticeably, this is one of the smallest galaxies for which we have a very accurate profile of the gravitating mass.

NGC3198 is a normal spiral about 2 times less luminous than the Milky Way. For a decade it held the record of the galaxy with the (HI) rotation curve showing the clearest evidence for Dark Matter.²⁰ Then, the record went to other galaxies with optical RCs, but recent radio measurements of a high-resolution¹⁸ has likely brought it back to this galaxy. In contrast with Orion, both its stellar and the HI disk are relevant.

The heart of this paper is that these two galaxies show without doubt a “Dark

Matter Phenomenon” but, when we analyse the issue in detail, we realise that well physically motivated halos of dark particles fail to account for their Rotation Curves. Our idea is to use these them to constrain proposed modifications of gravity: in the framework of those, can the baryonic matter alone account for the observed RCs when, in Standard Newtonian Gravity, the baryonic + dark matter together badly fail?

2. Newtonian limit of $f(R)$ gravity

The theory proposed by⁵ is an example of $f(R)$ theory of gravity.^{9,13} In these theories the gravitational action is defined to be:

$$\mathcal{S} = \int d^4x \sqrt{-g} [f(R) + \mathcal{L}_m], \quad (2)$$

where g is the metric determinant, R is the Ricci scalar and \mathcal{L}_m is the matter Lagrangian. They consider $f(R) = f_0 R^n$, where f_0 is a constant to give correct dimensions to the action and n is the slope parameter. The modified Einstein equation is obtained by varying the action with respect to the metric components.

Solving the vacuum field equations for a Schwarzschild-like metric in the Newtonian limit of weak gravitational fields and low velocities, the modified gravitational potential for the case of a point-like source of mass m , is given by Eq. (1), where the relation between the slope parameter n and β is given by:

$$\beta = \frac{12n^2 - 7n - 1 - \sqrt{36n^4 + 12n^3 - 83n^2 + 50n + 1}}{6n^2 - 4n + 2}. \quad (3)$$

Note that for $n = 1$ the usual Newtonian potential is recovered. The large and small scale behavior of the total potential constrain the parameter β to be $0 < \beta < 1$.

The solution Eq. (1) can be generalized to extended systems with a given density distribution $\rho(r)$ by simply writing:

$$\begin{aligned} \phi(r) &= -G \int d^3r' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} \left\{ 1 + \frac{1}{2} \left[\frac{|\mathbf{r} - \mathbf{r}'|^\beta}{r_c^\beta} - 1 \right] \right\} \\ &= \phi_N(r) + \phi_C(r), \end{aligned} \quad (4)$$

where $\phi_N(r)$ represents the usual Newtonian potential and $\phi_C(r)$ the additional correction. In this way, the Newtonian potential can be re-obtained when $\beta = 0$.

3. Data and methodology of the test

Let us remind following¹⁷⁻¹⁹ that for these galaxies we have high quality RC and a very good knowledge of the distribution of the luminous matter. Any result of the mass modelling could not be questioned on base of putative observational errors or biases. It is matter of fact that NFW halos + luminous matter model badly fits these very RCs^{17,19} and many others.^{1,14,16}

We decompose the total circular velocity into stellar and gaseous contributions. Available photometry and radio observations show that the stars and gas in these

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spirals are distributed in an infinitesimal thin and circular symmetric disk; from the HI flux we directly measure $\Sigma_{gas}(r)$ its surface density distribution (multiplied by 1.33 to take into account also the He contribution) In these galaxies, the stars follow the usual Freeman exponential thin disk:

$$\Sigma_D(r) = (M_D/2\pi R_D^2) e^{-r/R_D}. \quad (5)$$

M_D is the disk mass and it is kept as a free parameter, R_D is the scale length, measured directly from optical observations.

The distribution of the luminous matter has, to a good extent, a cylindrical symmetry and hence potential Eq. (4) reads

$$\phi(r) = -G \int_0^\infty dr' r' \Sigma(r') \int_0^{2\pi} \frac{d\theta}{|\mathbf{r} - \mathbf{r}'|} \left\{ 1 + \frac{1}{2} \left[\frac{|\mathbf{r} - \mathbf{r}'|^\beta}{r_c^\beta} - 1 \right] \right\}. \quad (6)$$

$\Sigma(r')$ is the surface density of the stars, given by Eq. (5), or of the gas, given by an interpolation of the HI measurements. β and r_c are, in principle, free parameters of the theory, with the latter perhaps galaxy dependent. We fix $\beta = 0.7$ to have agreement with previous results (see also⁸).

Defining $k^2 \equiv \frac{4r r'}{(r+r')^2}$, we can express the distance between two points in cylindrical coordinates as $|\mathbf{r} - \mathbf{r}'| = (r+r')^2(1 - k^2 \cos^2(\theta/2))$. The derivation of the circular velocity due to the marked term of Eq. (6), that we call $\phi_\beta(r)$, is now direct:

$$r \frac{d}{dr} \phi_\beta(r) = -2^{\beta-3} r_c^{-\beta} \pi \alpha (\beta - 1) G I(r), \quad (7)$$

where the integral is defined as

$$\mathcal{I}(r) \equiv \int_0^\infty dr' r' \frac{\beta - 1}{2} k^{3-\beta} \Sigma(r') \mathcal{F}(r), \quad (8)$$

with $\mathcal{F}(r)$ written in terms of confluent hyper-geometric function: $\mathcal{F}(r) \equiv 2(r+r') {}_2F_1[\frac{1}{2}, \frac{1-\beta}{2}, 1, k^2] + [(k^2 - 2)r' + k^2 r] {}_2F_1[\frac{3}{2}, \frac{3-\beta}{2}, 2, k^2]$.

The total circular velocity is the sum of each squared contribution:

$$V_{CT}^2(r) = V_{N,stars}^2 + V_{N,gas}^2 + V_{C,stars}^2 + V_{C,gas}^2, \quad (9)$$

where the N and C subscripts refer to the Newtonian and the additional modified potentials of the two different contributions (gas and stars) to the total potential Eq. (4).

In Fig. 1 the velocities are shown only in the ranges of r where their square are positive.

The RCs are χ^2 best-fitted with the free parameters: the scale length (r_c) of the theory and the gas mass fraction (f_{gas}) related to the disk mass simply by $M_D = M_{gas}(1 - f_{gas})/f_{gas}$ with the gas mass measured. The errors for the best fit values of the free parameters are calculated at one standard deviation.

Let us recall that we can write

$$V_{stars}^2(r) = (GM_D/2R_D) x^2 B(x/2), \quad (10)$$

where $x \equiv r/R_D$, G is the gravitational constant and the quantity $B = I_0 K_0 - I_1 K_1$ is a combination of Bessel functions.⁷

4. Results

We summarize the results of our analysis in Fig. 1. We find that the velocity model V_{CCT} is well fitting the RCs for very reasonable values of the stellar mass-to-light ratio. The resulting disk masses are $(3.7 \pm 0.8) \times 10^8 M_\odot$ and $(3.4 \pm 0.8) \times 10^{10} M_\odot$ respectively for Orion and NGC3198. The other parameters are: Orion $r_c = (0.013 \pm 0.002)$ kpc and gas fraction = $(55 \pm 20)\%$, NGC3198 $r_c = (0.4 \pm 0.05)$ kpc and gas

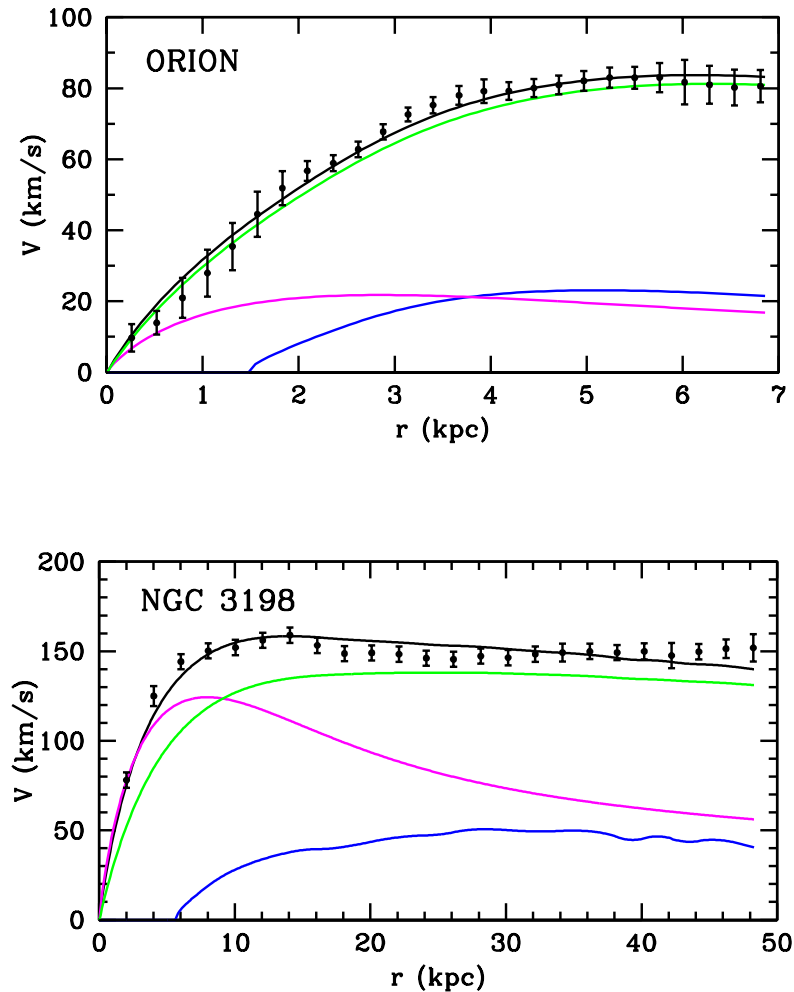


Fig. 1. Black: best-fit total circular velocity V_{CCT} . Blue: Newtonian gaseous contribution. Magenta: Newtonian stars contribution. Green: non-Newtonian gaseous and stars contributions to the model.

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fraction= $(29 \pm 10)\%$. The values of χ^2 are $\simeq 1$ confirming the success of the fit.

The value for the scale-length parameter r_c is found smaller for the less massive galaxy and larger for the more massive one, in line with previous results and with the idea of a scale dependent modification of gravity³.

5. Conclusions

Extended theories of gravity, created to tackle theoretical cosmological problems have something to say on another issue of Gravity, the Phenomenon of Dark Matter in galaxies. We have tested two objects with state of the art kinematical data that, in addition, are not accounted by the dark matter halo paradigm and we found that a scale dependent R^n Gravity is instead able to account for them. Extended theories of Gravity candidate themselves to explain the phenomenon of dark matter with only the luminous matter present in galaxies.

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References

1. F. Donato, G. Gentile, P. Salucci, Martins C. Frigerio, M. I. Wilkinson, G. Gilmore, E. K. Grebel, A. Koch and R. Wyse, *MNRAS* **397** (August 2009) 1169.
2. S. Capozziello, V. F. Cardone, S. Carloni and A. Troisi, *Phys.Lett.* **326** (June 2004) 292.
3. S. Capozziello, M. De Laurentis, *Annalen der Phys.* **524** (October 2012) 545.
4. S. Capozziello, V. F. Cardone and A. Troisi, *JCAP* **8** (August 2006) 001.
5. S. Capozziello, V. F. Cardone and A. Troisi, *MNRAS* **375** (March 2007) 1423.
6. S. Carroll, V. Duvvuri, M. Trodden and M. S. Turner, *Physical Review* **70** (August 2004) 043528.
7. K. C. Freeman, *ApJ* **160** (June 1970) 811.
8. C. Frigerio Martins and P. Salucci, *MNRAS* **381** (November 2007) 1103.
9. S. Capozziello and M. De Laurentis, *Physics Reports* **509** (December 2011) 167.
10. S. Nojiri and S.D. Odintsov, *Physics Reports* **505** (August 2011) 59.
11. T.P. Sotiriou and V. Faroni, *Reviews of Modern Physics* **82** (January 2010) 451.
12. M. Milgrom, *ApJ* **270** (July 1983) 365.
13. S. Nojiri and S. D. Odintsov, *Int.J.Geom.Meth.Mod.Phys.* **4** (January 2006) 115.
14. M. Persic and P. Salucci, F. Stel, *MNRAS* **281** (July 1996) 27.
15. V.C. Rubin, *Science* **220** (June 1983) 1339.
16. P. Salucci and M. Persic, *A&A* **351** (November 1999) 442.
17. N. Frusciante, P. Salucci, D. Vernieri, J. M. Cannon and E. C. Elson, *MNRAS* **426** (October 2013) 751.
18. G. Gentile, G.I.G. Jozsa, P. Serra, G.H. Heald, W.J.G. de Blok, F. Fraternali, M.T. Patterson, R.A.M. Walterbos and T. Osterloo, *A&A* **554** (June 2013) A125.
19. E.V. Karukes, G. Gentile and P. Salucci, in preparation.
20. T.S. van Albada, J.N. Bahcall, K. Begeman and R. Sancisi, *AJ* **295** (September 1985) 305.